

# How doubly-magic is the nucleus ${}^{78}_{28}\text{Ni}_{50}$ ?

B. Pfeiffer<sup>1,2</sup> and K.-L. Kratz<sup>1,2</sup>

<sup>1</sup>Institut für Kernchemie, Universität Mainz, Germany; <sup>2</sup>HGF VISTARS

Doubly-magic nuclei play an important role in nuclear structure theory as testbeds for shell model calculations. Data on neutron-rich nuclei in the regions of doubly-magic nuclei such as  ${}^{78}_{28}\text{Ni}_{50}$  and  ${}^{132}_{50}\text{Sn}_{82}$  have a decisive influence on nucleosynthesis calculations. These “longer-lived” waiting-point nuclei determine the duration of the r-process and the matter flow through the abundance maxima at the magic neutron numbers [1]. Remaining deficiencies prior to the abundance maxima in r-process calculations have been interpreted as signatures of new nuclear structure effects near the neutron drip-line, for example overestimation of the shell strength far from stability in global mass models such as FRDM and ETFSI-1. A

ferences of the two-neutron separation energies  $S_{2n}$  prior and behind a magic neutron number. Fig. 1 displays differences between experimental and theoretical  $S_{2n}$  values relative to the smoothly varying mass formula of Groote et al. [2] for Fe, Co and Ni isotopes.

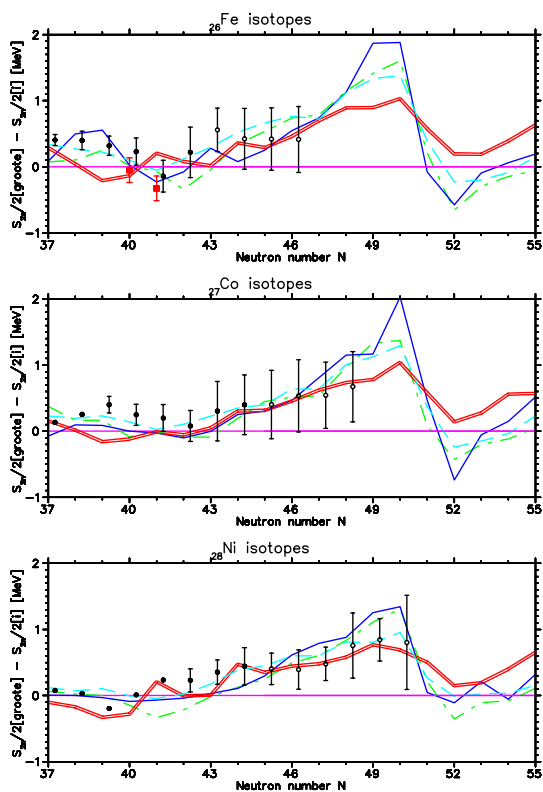


Figure 1: Deviations of experimental [black dots: 2003 mass evaluation [7], red squares: FRS/ESR measurements [8]] and theoretical  $S_{2n}$  values from the smoothly varying Groote mass formula [2] across the shell gap at  $N=50$  are shown for  ${}_{26}\text{Fe}$  (upper part),  ${}_{27}\text{Co}$  (middle part) and  ${}_{28}\text{Ni}$  (lower part) isotopes. [Theoretical masses: Groote: magenta, FRDM: red [3], ETFSI-1: cyan [4], HFB-2: green [5], HFB-8: blue [6]]

weakening (“quenching”) of spherical shells with increasing isospin, resulting in a gradual setting in of collectivity, has been predicted by recent HFB calculations, and is well established for the lower neutron-magic numbers  $N=20$  and  $N=28$ .

In this context, it is of interest to determine the mutual influence of the proton and neutron magic numbers far from stability. The shell strength can be derived from the dif-

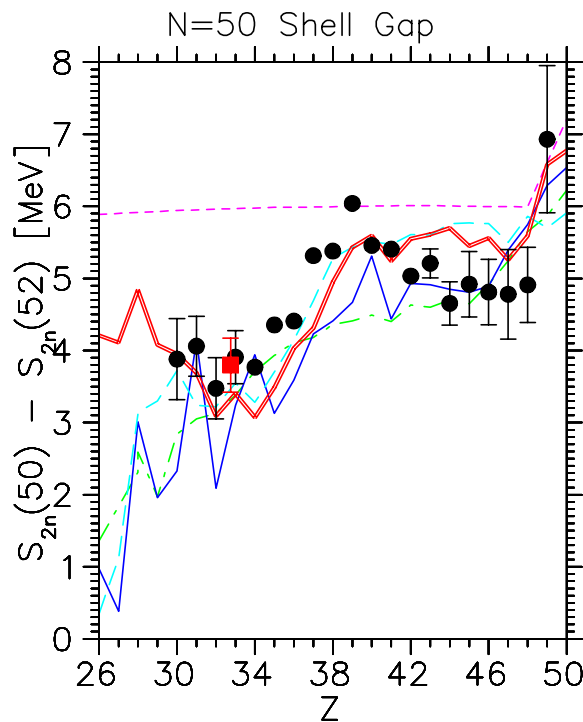


Figure 2: The  $N=50$  shell gap as a function of  $Z$ . Experimental and extrapolated values from the 2003 mass evaluation [7] (black circles) and experimental values from direct mass measurements at ESR/FSR [red square] are compared to theoretical mass models. [Same colour coding for the mass models as in Fig. 1.]

Fig. 2 shows the  $N=50$  shell gap over a wide  $Z$  range. Most global mass models predict a local maximum for doubly-magic  ${}^{78}\text{Ni}$ . So far, below  $Z=30$  no experimental masses have been determined, so that this prediction cannot be verified. With upgraded U-beams, direct mass measurements at the FRS-ESR of GSI will in future extend the range of experimental masses in this region to more neutron-rich isotopes.

## References

- [1] K.-L. Kratz et al., *Ap. J.* **403**, 216 (1993).
- [2] H. von Groote et al., *ADNDT* **17**, 418 (1976).
- [3] P. Möller et al., *ADNDT* **66**, 131 (1997).
- [4] Y. Aboussir et al., *ADNDT* **61**, 127 (1995).
- [5] S. Goriely et al., *ADNDT* **77**, 311 (2001).
- [6] M. Samyn et al., *Phys. Rev.* **70**, 044309 (2004).
- [7] G. Audi et al., *Nucl. Phys.* **A729**, 3 (2003).
- [8] M. Matos, PhD thesis, Giessen (2004).