

Influence of thermally excited isomers on r-process abundances?

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The influence of nuclear physics input as well as astrophysical parameters on the abundances of r-process nucleosynthesis can be studied systematically within the “waiting-point” concept [1]. Calculations applying theoretical nuclear masses and $T_{1/2}$ and P_n values from the FRDM or the ETFSI-Q mass models allow a satisfactory reproduction of the global isotopic abundances ($N_{r,\odot}$). Nevertheless, there still exist deviations in limited mass regions as the left wing of the $A \approx 130$ peak (see black line in Fig. 1). Even replacing the theoretical values with experimental ones (as far as available) does not solve the problem (see Ref. [3]).

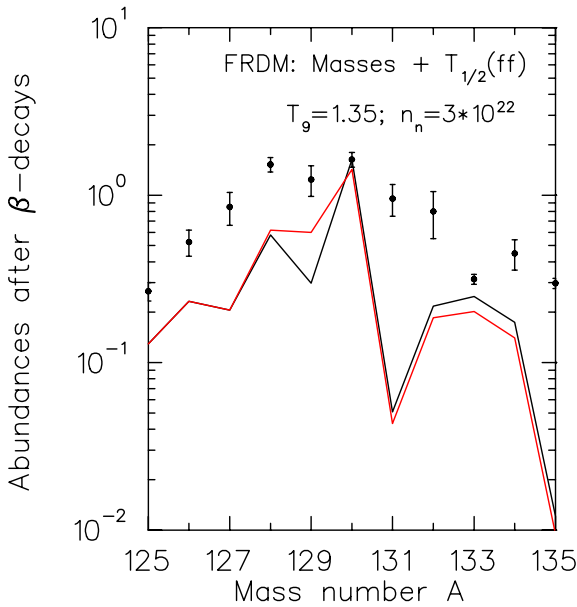


Figure 1: Solar r-process abundances calculated with masses and half-lives from the FRDM mass model [2]. The half-life of ^{129}Ag is either taken as the $\pi g_{9/2}$ g.-s. [black line] or as $T_{1/2}=180$ ms [red line] (see text).

As the abundances are proportional to the half-lives of the precursor nuclei in the r-process boulevard, longer half-lives for the $N=82$ nuclei with $Z \leq 47$ could be a solution. The r-process nucleosynthesis takes place in a heat bath with temperatures in excess of $1 T_9$. In the case of ^{129}Ag we have observed indications to the existence of an isomeric β -decaying $\pi p_{1/2}$ state with $T_{1/2}=160$ ms in addition to the $\pi g_{9/2}$ ground-state (g.-s.) decay with $T_{1/2}=46$ ms [4]. The excitation energy of the isomeric state is not known experimentally. QRPA calculations yield 500 keV. The stellar half-life of ^{129}Ag is shown in Fig. 2 as a function of the equilibrium temperature in units of 10^9 K (T_9). At the freeze-out temperature of $1.35 T_9$, the stellar half-life is only marginally enhanced due to the high excitation energy, not influencing the abundance at $A=129$.

If one follows the γ -decay of unbound states populated

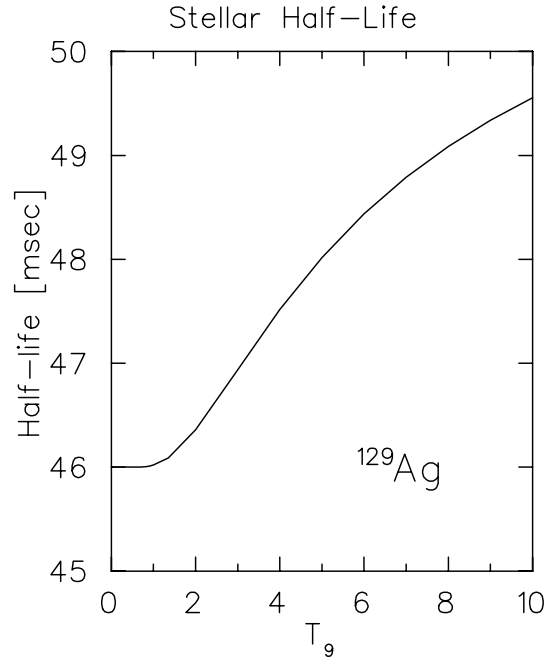


Figure 2: Stellar half-life of ^{129}Ag as a function of temperature. $T_{1/2}$ for the $\pi g_{9/2}$ g.-s. is 46 ms and 160 ms for the $\pi p_{1/2}$ β -isomeric state, respectively. The excitation energy of the isomeric state is assumed as 500 keV according to QRPA calculations.

in neutron capture on ^{128}Ag under non-equilibrium conditions, about one third of the rays populate the isomeric state [5]. From this a stellar half-life of about 80 ms can be deduced, much longer than in the case of the thermal excitation under equilibrium condition. The isotopic abundances obtained in replacing the short g.-s. half-life of ^{129}Ag with this longer value are displayed in Fig. 1 [red line]. The abundance for $A=129$ is now closer to the observed value.

The general solution for the calculated underabundances in the left wing of the $A=130$ peak has to be sought in nuclear structure different from the assumptions underlying the global mass models. Applying single-particle energies which reproduce the surprisingly high-lying 1^+ state in ^{130}Cd in Nilsson model calculations of β -decay half-lives of nuclides in the left wing, longer half-lives were obtained [3].

References

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